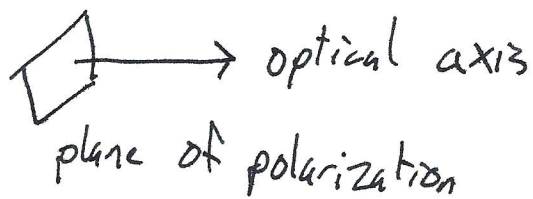


Phys 110B 30 Nov 20

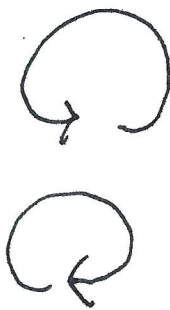
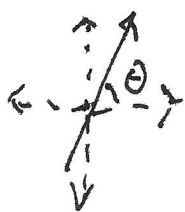
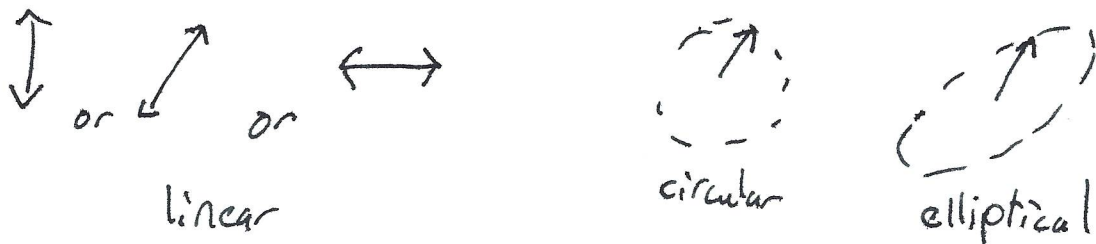
Polarization "Jones Calculus"

well-collimated beam of EM radiation



polarization \rightarrow direction of \vec{E} field

polarization states \leftrightarrow 2D complex vectors



Jones vector complex 2D polarization vector in some basis

{ Horizontal $\leftrightarrow \vec{H} \quad |H\rangle$
 { vertical $\updownarrow \vec{V} \quad |V\rangle$

{ diagonal $\nearrow \vec{D} \quad |D\rangle$
 { anti-diagonal $\searrow \vec{A} \quad |A\rangle$

{ left circular $\curvearrowleft \vec{L} \quad |L\rangle$
 { right circular $\curvearrowright \vec{R} \quad |R\rangle$

Orthonormal basis makes life easier

$$\vec{H}^\dagger \vec{H} = 1 \quad \vec{V}^\dagger \vec{V} = 1 \quad \vec{H}^\dagger \vec{V} = 0$$

$$\langle H|H\rangle = 1 \quad \langle V|V\rangle = 1 \quad \langle H|V\rangle = 0$$

In this \vec{H}, \vec{V} basis

$$\vec{H} \cong \begin{bmatrix} 1 \\ 0 \end{bmatrix} \quad \vec{V} \cong \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\vec{E} = \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = \alpha \begin{bmatrix} 1 \\ 0 \end{bmatrix} + \beta \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\langle \vec{S} \rangle = \sqrt{\frac{E}{\mu}} \frac{c}{8\pi} (\langle E_x^\dagger E_x \rangle + \langle E_y^\dagger E_y \rangle) \propto |\alpha|^2 + |\beta|^2 \propto \vec{E}^\dagger \vec{E}$$

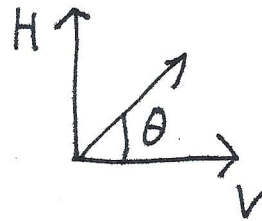
In \vec{H}, \vec{V} basis

$$\vec{D} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \quad \vec{A} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

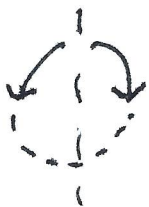
georgi waves convention
for \vec{L}, \vec{R}

$$\vec{L} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix} \quad \vec{R} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$$

$$\hat{E}_\theta = \begin{bmatrix} \cos\theta \\ \sin\theta \end{bmatrix} \quad \hat{E}_{\theta+\frac{\pi}{2}} = \begin{bmatrix} \sin\theta \\ -\cos\theta \end{bmatrix}$$

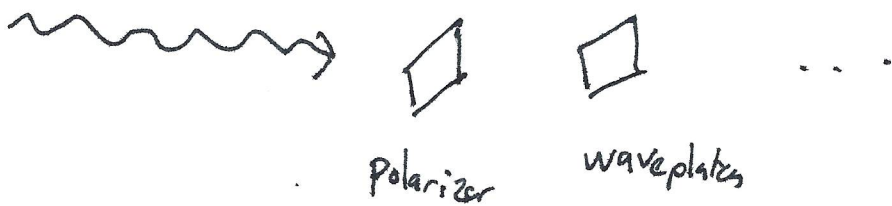


in \vec{L}, \vec{R} basis



$$\frac{\vec{L} + \vec{R}}{\sqrt{2}} = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix} + \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix} \right) = \frac{1}{2} \begin{bmatrix} 2 \\ 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \vec{H}$$

$$\frac{\vec{L} - \vec{R}}{\sqrt{2}i} = \frac{1}{\sqrt{2}} \frac{1}{\sqrt{2}} \frac{1}{i} \begin{bmatrix} 0 \\ 2i \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} = \vec{V}$$



"device" in the beam line \leftrightarrow 2×2 complex matrix

Linear Polarizer (idealized case)

$$P_{\vec{H}} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \end{bmatrix} = \vec{H} \vec{H}^\dagger$$

horizontal

$$P_{\vec{V}} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \end{bmatrix} = \vec{V} \vec{V}^\dagger$$

$$P_{\theta} = \begin{bmatrix} \cos\theta \\ \sin\theta \end{bmatrix} \begin{bmatrix} \cos\theta & \sin\theta \end{bmatrix} = \begin{bmatrix} \cos^2\theta & \cos\theta \sin\theta \\ \sin\theta \cos\theta & \sin^2\theta \end{bmatrix}$$

P_{θ} = Hermitian Projections

$$P_{\theta}^T = P_{\theta} \quad P_{\theta} P_{\theta} = P_{\theta} \quad (\text{idempotent})$$

$$(\hat{u} \hat{u}^T) = \hat{u}^T \hat{u}^T = \hat{u} \hat{u}^T$$

$$(\hat{u} \hat{u}^T)(\hat{u} \hat{u}^T) = \hat{u} (\hat{u}^T \hat{u}) \hat{u}^T = \hat{u} \hat{u}^T$$

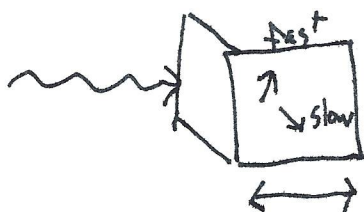
$$\hat{u} \cdot \hat{u} = 1$$

filter

$$F = f \mathbb{1} \quad 0 \leq f \leq 1$$

↑
identity

Waveplates birefringence



Quarter-wave plate

$$\Delta_H = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}_{H,V} = P_H + iP_V$$

↑
fast axis

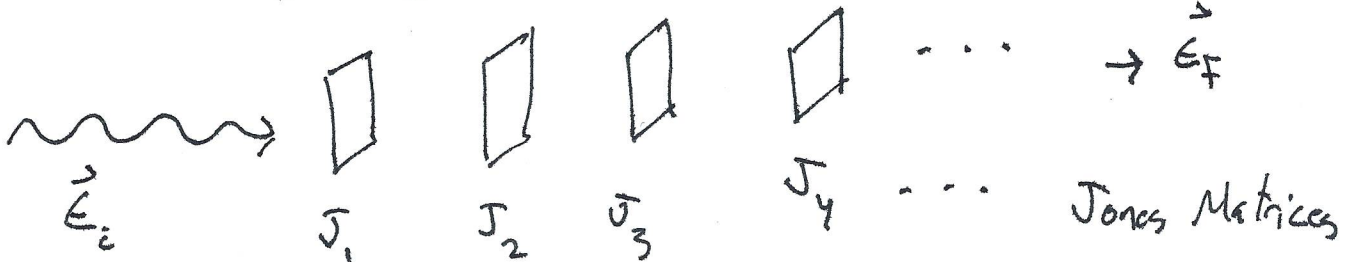
$$\Delta_{\theta} = P_{\theta} + iP_{\theta + \frac{\pi}{2}}$$

↑
fast axis

Half waveplate

$$A_{\theta} = A_0 A_{\theta}$$

$$H_H = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$



$$\vec{E}_f = J_n \cdots J_4 J_3 J_2 J_1 \vec{E}_i$$

$$\vec{E}_f \propto \langle \vec{s} \rangle$$