

Phys 110B 18 Nov 20

Radiation Reaction

non-relativistic case

$$\vec{F}_{\text{rad}} \cdot \vec{v} = -P_{\text{Larmor}} = \frac{-\mu_0 q^2 a^2}{6\pi c}$$

$$\int_{t_1}^{t_2} \vec{F}_{\text{rad}} \cdot \vec{v} dt = \frac{-\mu_0 q^2}{6\pi c} \left\{ \vec{v} \cdot \frac{d\vec{v}}{dt} \Big|_{t_1}^{t_2} - \int_{t_1}^{t_2} \vec{v} \cdot \frac{d^2\vec{v}}{dt^2} dt \right\}$$

||
0

suggests

$$\vec{F}_{\text{rad}} = \frac{\mu_0 q^2}{6\pi c} \dot{\vec{a}} \quad \text{Abraham-Lorentz Force}$$

$$\vec{F} = m\vec{a} \rightarrow \vec{F}_{\text{ext}} + \vec{F}_{\text{rad}} = m\vec{a} = m\ddot{\vec{x}}$$

$$\vec{F}_{\text{ext}}(\vec{x}, \vec{v}) \quad \vec{F}_{\text{rad}} \propto \dot{\vec{a}} = \ddot{\vec{x}}$$

changes EOM to a 3rd order differential equation

in particular suppose $\vec{F}_{\text{ext}} = \vec{0}$ $\vec{x} = \dot{\vec{x}} = \dot{\vec{a}} = \vec{0} \quad t < 0$

$$\vec{a}(t) = \begin{cases} \vec{0} \\ \vec{a}_0 e^{t/\tau} \end{cases}$$

$$\tau = \frac{1}{4\pi\epsilon_0} \frac{2}{3} \frac{e^2}{m c^3} = \frac{2}{3} \frac{e^2}{4\pi\epsilon_0 m c^3}$$

q_0 vs $q_0 \dots$ 

pre "runaway acceleration" problem

Can eliminate the "runaway" solutions by imposing

$$\lim_{t \rightarrow \infty} \vec{a}(t) = \vec{0} \text{ if } \vec{F}_{\text{ext}} \text{ turns off}$$

impose via turning ODE into an integral equation

$$m \dot{v}(t) = \int_0^{\infty} e^{-\tau s} \vec{F}(t + \tau s) ds$$

avoids runaways but exhibits "pre-acceleration"
acceleration now depends on acceleration in future

Look at Relativistic case:

$$\frac{d}{d\tau} p^\mu = e F^{\mu\nu} U_\nu = K_{\text{ext}}^\mu + K_{\text{rad}}^\mu$$

$$\frac{d}{d\tau} p^\mu + K_{\text{ext}}^\mu + U^\mu \left[\frac{2}{3} \frac{e^2}{c^3 4\pi\epsilon_0} \ddot{U}^\alpha \dot{U}_\alpha \right] = \frac{d}{d\tau} \text{ Liénard}$$

that "balances" the power loss in time-integrated sense
(if we ignore boundary terms)

Inconsistent as a 4-force

$$U^\mu K_\mu \neq 0 \text{ in general}$$

Lorentz - Abraham - Dirac Force

only possibility that depends on $u^\mu, \dot{u}^\mu, \ddot{u}^\mu$

$$\dot{p}^\mu = K_{\text{ext}}^\mu + K_{\text{self}}^\mu$$

$$K_{\text{self}}^\mu = \frac{2}{3} \frac{e^2}{4\pi\epsilon_0} \left[\ddot{u}^\mu - \frac{u^\mu}{c^2} \dot{u}^\alpha \dot{u}_\alpha \right]$$

$$K_{\text{self}}^\mu u_\mu = 0$$

$$K_{\text{self}}^\mu = K_{\text{Schatt}}^\mu + K_{\text{rad}}^\mu \downarrow$$

$$\downarrow \frac{2}{3} \frac{e^2}{4\pi\epsilon_0} \ddot{u}^\mu$$

$$-\frac{2}{3} \frac{e^2}{4\pi\epsilon_0} u^\mu \dot{u}^\alpha \dot{u}_\alpha$$

total time derivative
reversible
involve near fields

not a total time derivative
irreversible
involves exchange of
energy with radiation fields

$$F_{\text{ext}}^{\mu\nu} = F_{\text{incident}}^{\mu\nu} + \int \mathcal{L}_{\text{rad}}^{\mu\nu}$$

$$A^\mu = A^\mu_{\text{incident}} + \int G_r J^\mu$$

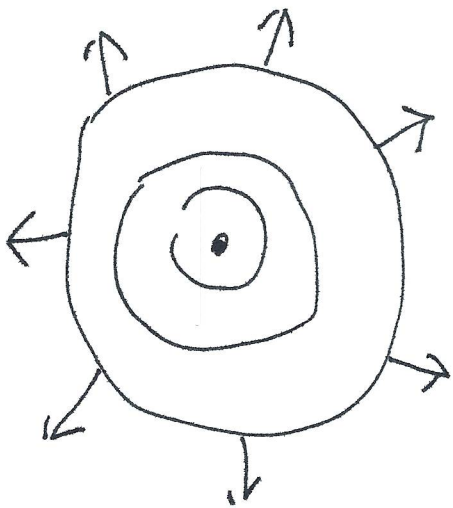
$$A^\mu = A^\mu_{\text{out}} + \int G_a J^\mu$$

$$A^\mu_{\text{out}} - A^\mu_{\text{in}} = \int d^3x' (G_r(\vec{x} - \vec{x}') - G_a(\vec{x} - \vec{x}')) J^\mu(\vec{x}')$$

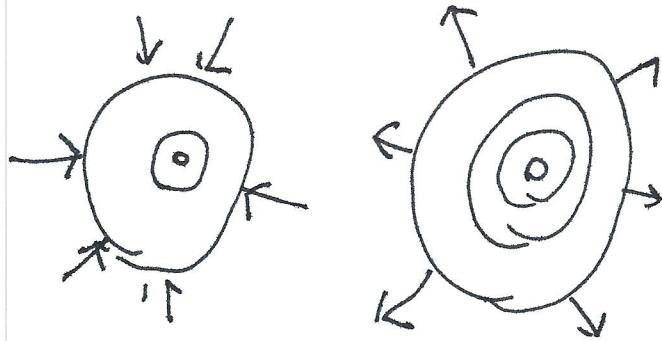
↑ satisfy the source-free Maxwell's equations everywhere

$$A_{\text{out}} - A_{\text{in}} = A_{\text{rad}}$$

"radiation" fields Dirac



outgoing
radiation



Dirac Radiation

A_{rad} is regular everywhere

(does not blow up at the source)

$$F^{\mu\nu} = F_{ret}^{\mu\nu} = \frac{1}{2}(F_{ret}^{\mu\nu} - F_{adv}^{\mu\nu}) + \frac{1}{2}(F_{ret}^{\mu\nu} + F_{adv}^{\mu\nu})$$

actual fields (near + far) for all sources

$$= \frac{1}{2}F_{rad}^{\mu\nu} + \bar{F}^{\mu\nu}$$

↑ regular everywhere

↑ blows up at point charge sources

$$\dot{p}^M = K_{ext}^M + K_{self}^M - \delta m \dot{u}^M$$

" $m \dot{u}^M$ " $\int F_{ext}^{\mu\nu} u_\nu$ ↑ infinite constant for a point charge

$$(m + \delta m) \dot{u}^M = K_{ext}^M + K_{self}^M$$

↑ from $\vec{F}^{\mu\nu}$

external Lorentz force

↑ from $F_{ext}^{\mu\nu}$

Abraham-Lorentz-Dirac self force

↑ $F_{rad}^{\mu\nu}$

$(m_{base} + \delta m) = m_{exp}$ observed mass

without radiation reaction effects

Renormalization of particle mass

$$m_{\text{base}} + \delta m = m_{\text{exp}} \text{ is to be finite}$$

\uparrow \downarrow
 $\pm \infty$ $\pm \infty$ for a point charge

To avoid pre-acceleration/runaway

Landau-Lifshitz - Spohn

$$m \dot{u}^\mu = 2 F_{\text{ext}}^{\mu\nu} u_\nu + \frac{2}{3} \frac{e^2}{c^3 4\pi\epsilon_0} \left(\frac{e}{m} u^\sigma \partial_\sigma F^{\mu\nu} u_\nu \right. \\ \left. + \left(\frac{e}{m}\right)^2 F^{\mu\alpha} F_{\alpha\nu} u_\nu + \left(\frac{e}{m}\right)^2 \frac{1}{c^2} F^{\sigma\alpha} F_{\alpha\nu} u_\sigma u_\nu u^\mu \right)$$

closest thing we have to a self-consistent form EOM
for a classical point charge