

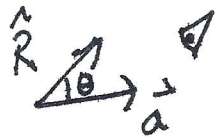
Phys 110B 09 Nov 20

$$\vec{E}_{\text{rad}} = \frac{q}{4\pi\epsilon_0} \left[\frac{\hat{R} \times (\hat{R} - \vec{\beta}) \times \dot{\vec{\beta}}}{(1 - \vec{\beta} \cdot \hat{R})^3 c R} \right]_{\text{ret}}$$

$$\vec{B}_{\text{rad}} = \left[\frac{\hat{R} \times \vec{E}}{c} \right]_{\text{ret}}$$

$$\vec{S}_{\text{rad}} = \frac{1}{\mu_0 c} |\vec{E}_{\text{rad}}|^2 \hat{R}$$

Particle is instantaneously at rest



$\vec{\beta} = \vec{0}$ instantaneously

$$R^2 \vec{S}_{\text{rad}} \cdot \hat{R} = \frac{dP}{d\Omega} = \dots = \frac{\mu_0 q^2}{16\pi^2 c} a^2 \sin^2 \theta$$

$$P = \int \frac{dP}{d\Omega} d\Omega = \frac{\mu_0 q^2}{6\pi c} a^2 \quad \text{Larmor Formula}$$

$$d\Omega = \sin\theta d\theta d\phi \quad \int \frac{1}{4\pi\epsilon_0} \frac{2}{3} q^2 a^2$$

Alternate derivation

Recall for an electric dipole $\omega^2 \vec{p} \propto \ddot{\vec{p}}$

far-field

$$\vec{E} = \frac{\mu_0}{4\pi} \frac{1}{r} \left[\hat{r} \times (\hat{r} \times [\ddot{\vec{p}}]) \right] \propto \frac{\mu_0}{4\pi} [\ddot{\vec{p}}] \frac{\sin\theta}{r} \hat{\theta}$$

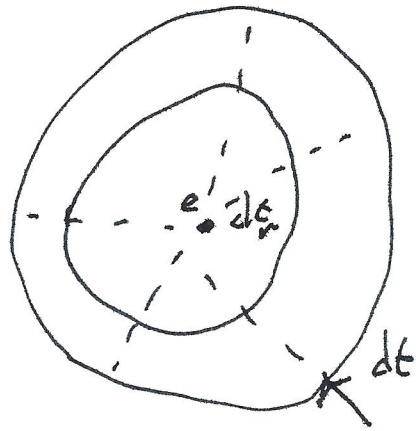
$$\vec{B} = \frac{-\mu_0}{4\pi c} \frac{1}{r} \left[\hat{r} \times [\dot{\vec{p}}] \right] \propto \frac{\mu_0}{4\pi} [\dot{\vec{p}}] \frac{\sin\theta}{r} \hat{\phi}$$

$$\vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B}) = \frac{\mu_0}{16\pi^2 \epsilon_0} \left[\frac{d^2 \vec{p}}{dt^2} \right] \frac{\sin^2 \theta}{r^2} \hat{r}$$

$$\Rightarrow P = \frac{\mu_0}{6\pi \epsilon_0} [\dot{p}']^2 \quad p \sim qd$$

$$\dot{p} \sim q \dot{d} \sim qa$$

$$P = \frac{\mu_0 q^2}{6\pi \epsilon_0} a^2 \quad \text{Larmor Formula}$$



(Relativistic) Radiated Power

$$\frac{dP}{d\Omega} = R^2 \frac{d\vec{S}}{d\Omega} \cdot \hat{R} \cdot \frac{dt}{dt_r} = R^2 \frac{d\vec{S}}{d\Omega} \cdot \hat{R} [1 - \vec{\beta} \cdot \hat{R}] \quad \text{if } \vec{\beta} = \vec{0}$$

then $dt_r = dt$

Power observed per solid angle if $\vec{\beta} \neq \vec{0}$ then $k \neq 0$

Energy per observer time per solid angle

Not a natural measure of particle's radiation

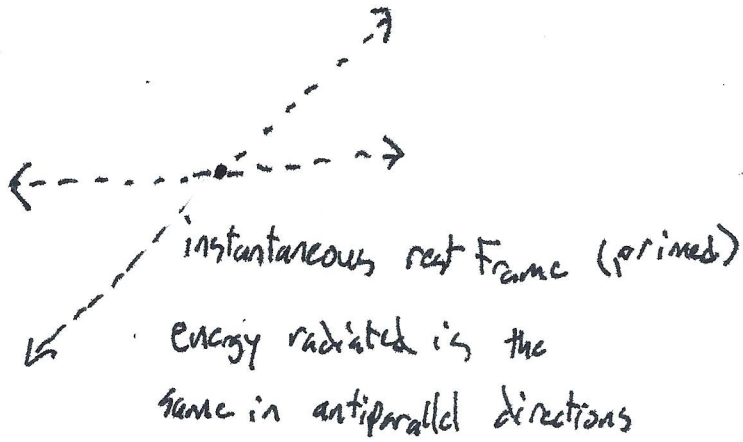
$$\frac{dP(t_r)}{d\Omega} = \dots = \frac{q^2}{16\pi^2 \epsilon_0 c} \left[\frac{\left| \hat{R} \times ((\hat{R} - \vec{\beta}) \times \dot{\vec{\beta}}) \right|^2}{(1 - \hat{R} \cdot \vec{\beta})^5} \right]$$

$$P(t_r) = \dots = \frac{\mu_0 q^2}{6\pi \epsilon_0} [\gamma^6] \left[|\dot{\vec{a}}|^2 - |\vec{\beta} \times \dot{\vec{a}}|^2 \right] \geq 0$$

Relativistic generalization of Larmor Formula

Liénard's Formula

* Emitted power (radiated by a point charge) must be a Lorentz Invariant



$$\begin{bmatrix} \frac{\Delta E'}{c} \\ \Delta p'_{||} \\ \Delta p'_{\perp} \end{bmatrix} \leftarrow \begin{array}{l} \text{energy radiated in a time } \Delta t' \\ \text{4-vector} \end{array}$$

$$\begin{bmatrix} \gamma & \beta\gamma & 0 \\ \beta\gamma & \gamma & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \frac{\Delta E'}{c} \\ \Delta p'_{||} \\ \Delta p'_{\perp} \end{bmatrix} = \begin{bmatrix} \frac{\Delta E}{c} \\ \Delta p_{||} \\ \Delta p_{\perp} \end{bmatrix}$$

this term goes away as it averages to zero

$$\Delta E = \gamma (\Delta E' + \beta c \Delta p'_{||})$$

$$\Delta E = \gamma \Delta E'$$

lab Frame proper frame of source

back to lab frame

$$\Delta t' = \Delta \tau = \frac{\Delta t}{\gamma} \quad \Delta t = \gamma \Delta t'$$

$$\frac{\frac{\Delta E}{\Delta t}}{\Delta t} = \frac{\Delta E}{\Delta t} = \frac{\gamma \Delta E'}{\gamma \Delta t'} = \frac{\Delta E'}{\Delta t'}$$

emitted power in lab frame emitted power in rest frame

$\therefore \frac{\Delta E}{\Delta t}$ is a Lorentz scalar

Quantum Mechanically, the number of emitted photons per proper time interval must be invariant.

1) Emitted power is a Lorentz scalar

2) $e, m, q, P^\mu, \frac{dP^\mu}{d\tau}$, etc.
of
source

3) reduce to Larmor Formula as $|\vec{\beta}| \rightarrow 0$

the only possibility

$$P = \sum_{\mu} \frac{dP^\mu}{d\tau} \frac{dP_\mu}{d\tau}$$

↑
scalar depending

on e, m, q, \dots

agreement with Larmor $\rightarrow \sum_{\mu} \frac{dP^\mu}{d\tau} \frac{dP_\mu}{d\tau} = \frac{\mu_0 q^2}{6\pi m^2 c}$

Agrees with earlier expression for Liénard Formula

radiated energy-momentum ΔP^μ must be a 4-vector

$$\frac{d}{d\tau} \Delta P^\mu = \frac{1}{4\pi\epsilon_0} \frac{2}{3} \frac{q^2}{c^3} a^\nu a_\nu U^\mu$$

↑
emitter's proper time

↑
4-acceleration
of particle

←
4-velocity of source

$$P = \frac{\mu_0 e^2}{6\pi c} [\gamma^6] \left[|\dot{\vec{a}}|^2 - |\dot{\vec{\beta}} \times \dot{\vec{a}}|^2 \right]$$

Case 1: Bremsstrahlung - "braking radiation"

$\dot{\vec{\beta}}, \dot{\vec{a}}$ colinear (parallel or antiparallel)

