

Griffiths 10.6

a) Show it's always possible to choose $\vec{\nabla} \cdot \vec{A} + \mu_0 \epsilon_0 \partial_t \phi = 0$

b) Show that $\phi = 0$ is always possible

c) Argue that the Coulomb gauge is always possible

a) Suppose $\vec{\nabla} \cdot \vec{A} \neq \mu_0 \epsilon_0 \partial_t \phi$

$$\vec{\nabla} \cdot \vec{A} + \mu_0 \epsilon_0 \partial_t \phi = \psi$$

Consider gauge transformation generalized by λ

$$\vec{\nabla} \cdot \vec{A}' + \mu_0 \epsilon_0 \partial_t \phi' = 0 = \vec{\nabla} \cdot \vec{A} + \nabla^2 \lambda + \mu_0 \epsilon_0 \partial_t \phi - \mu_0 \epsilon_0 \partial_t^2 \phi$$

$$\Rightarrow (\nabla^2 + \mu_0 \epsilon_0 \partial_t^2) \lambda = \psi$$

↑
forced wave equation

✓ L-L gauge

$$\square \chi = 0$$

b) \vec{A}, ϕ

$$\chi(\vec{x}, t) = \int_0^t dt' \phi(\vec{x}, t')$$

$$\phi' = \phi - \partial_t \chi = 0$$

✓ Temporal Gauge

$$c) \vec{\nabla} \cdot (\vec{A} + \vec{\nabla} \chi) = 0$$

$$\nabla^2 \chi = -\vec{\nabla} \cdot \vec{A} \quad \checkmark \quad \text{Coulomb Gauge}$$

$$d) \vec{A} = \vec{0} \quad \leftarrow 3 \text{ conditions}$$

$$\Rightarrow \vec{\nabla} \times \vec{A} = \vec{0} \Rightarrow \vec{B} = \vec{0} \quad X$$

gauge \rightarrow 1 condition

Griffiths 10.7th

$$\rho(\vec{x}, t) = \delta^3(\vec{x}) q(t) \quad \vec{r} = \vec{x}$$

$$\vec{J}(\vec{x}, t) = -\frac{1}{4\pi} \dot{q}(t) \frac{\hat{r}}{r^2}$$

a) verify charge conservation holds

b) Find a ϕ, \vec{A} in the Coulomb Gauge

c) Find the corresponding \vec{E}, \vec{B} fields

d) verify they satisfy Maxwell's Equations

$$a) \vec{\nabla} \cdot \vec{J}(\vec{x}, t) = -\frac{\dot{q}}{4\pi} \vec{\nabla} \cdot \frac{\hat{r}}{r^2} = -\frac{\dot{q}}{4\pi} 4\pi \delta^3(\vec{r}) = -\dot{q} \delta^3(\vec{r}) = -\frac{\partial \rho}{\partial t}$$

$$b) \phi(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t)}{|\vec{r} - \vec{r}'|} dV = \frac{q(t)}{4\pi\epsilon_0} \int \frac{\delta^3(\vec{r}')}{|\vec{r} - \vec{r}'|} dV = \frac{q(t)}{4\pi\epsilon_0 r}$$

$$\text{By symmetry, } \vec{B} = \vec{0} \Rightarrow \vec{\nabla} \times \vec{A} = \vec{0}$$

$$\vec{\nabla} \cdot \vec{A} = 0$$

$$\vec{A} \rightarrow \vec{0} \text{ as } |\vec{r}| \rightarrow \infty$$

$$\vec{A}(\vec{r}, t) = \vec{0}$$

$$c) \vec{E} = -\vec{\nabla}\phi - \partial_t \vec{A} = -\vec{\nabla}\phi = \frac{1}{4\pi\epsilon_0} \frac{q(t)}{r^2} \hat{r}$$

$$\vec{B} = \vec{0}$$

$$d) \vec{\nabla} \cdot \vec{E} = \frac{q(t)}{4\pi\epsilon_0} \vec{\nabla} \cdot \left(\frac{\hat{r}}{r^2} \right) = \frac{q(t)}{4\pi\epsilon_0} 4\pi\delta^3(\vec{r}) = \frac{q(t)}{\epsilon_0} \delta(\vec{r}) = \rho(\vec{r}, t) \quad \checkmark$$

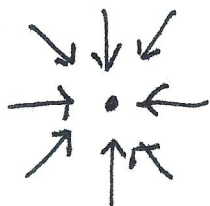
$$\vec{\nabla} \cdot \vec{B} = 0 \quad \checkmark$$

$$\vec{\nabla} \times \vec{E} = \vec{0} = -\frac{\partial \vec{B}}{\partial t} \quad \checkmark$$

by symmetry radial symmetry has no curl

$$\vec{0} = \vec{\nabla} \times \vec{B} = \mu_0 \vec{\nabla} \times \vec{J} + \mu_0 \epsilon_0 \partial_t \vec{E}$$

$$= \mu_0 \left(-\frac{1}{4\pi} \frac{\dot{q}}{r^2} \hat{r} \right) + \mu_0 \epsilon_0 \left(\frac{\dot{q}}{4\pi\epsilon_0} \frac{\hat{r}}{r^2} \right) = \vec{0} \quad \checkmark$$



Griffiths 10.13

Suppose $\vec{J}(\vec{r})$ is independent of time

a) Find the most general form of $\rho(\vec{r}, t)$

b) Find $\vec{E}(\vec{r}, t)$

$$a) \vec{\nabla} \cdot \vec{J} + \partial_t \rho = 0$$

$$\partial_t \rho = -\vec{\nabla} \cdot \vec{J} = \text{independent of time, } \rho \text{ not necessarily time independent}$$

$$\rho(\vec{r}, t) = \rho(\vec{r}, 0) + \dot{\rho}(\vec{r}, 0)t$$

ρ can at most be linear in time

$$b) \vec{E}(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \left(\frac{(\rho(\vec{r}', t'))_{ret}}{R^2} + \frac{(\partial_{t'} \rho(\vec{r}', t'))_{ret}}{cR} + \frac{(\partial_{t'}^2 \vec{J})_{ret}}{c^2 R} \right) d^3\vec{r}'$$

$$= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \left(\frac{\rho(\vec{r}', 0) + t_r \dot{\rho}(\vec{r}', 0)}{R^2} + \frac{\dot{\rho}(\vec{r}', 0)}{cR} \right) \hat{R}$$

$\downarrow 0$

$$t_r = t - \frac{R}{c}$$

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \left(\frac{\rho(\vec{r}', 0)}{R^2} + \frac{\dot{\rho}(\vec{r}', 0)}{R^2} - \frac{\dot{\rho}(\vec{r}', 0) R}{R^2 c} + \frac{\dot{\rho}(\vec{r}', 0)}{cR} \right)$$

$$= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \left(\frac{\rho(\vec{r}', 0) + \dot{\rho}(\vec{r}', 0) t}{R^2} \right) \hat{R}$$

$$= \frac{1}{4\pi\epsilon_0} \int d^3\vec{r}' \left(\frac{\rho(\vec{r}', t)}{R^2} \right) \hat{R}$$

Coulomb Field due to instantaneous $\rho(\vec{r}, t)$